Mini-Symposium on Topological Quantum Computation

National University of Ireland at Maynooth

# **D-Colexes & Topological Color Codes**

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## **References**

#### 2-Colexes

H.Bombin, M.A. Martin-Delgado, "Topological Quantum Distillation", Phys. Rev. Lett. **97** 180501 (2006)

#### **3-Colexes**

H.Bombin, M.A. Martin-Delgado, "Topological Computation without Braiding", quant-ph/0605138

#### **Topological order**

H.Bombin, M.A. Martin-Delgado, "Exact Topological Quantum Order in D=3 and Beyond: Branyons and Brane-Net Condensates", Phys. Rev. B accepted

## **Outline**

#### Stabilizer codes

- Transversal gates. Reed-Muller codes and universality.

#### Topological stabilizer codes

- Surface codes.

#### • 2-Colexes

- 2D-lattice. Stabilizer.
- Strings and string-nets.
- Implementation of the Clifford group.

#### 3-Colexes

- Universal quantum computation.

• A **stabilizer code**<sup>1</sup> *C* of length *n* is a subspace of the Hilbert space of a set of *n* qubits. It is defined by a stabilizer group *S* of Pauli operators, i.e., tensor products of Pauli matrices.

 $|\psi\rangle\in\mathcal{C}\qquad\iff\qquad\forall\,s\in\mathcal{S}\quad s|\psi\rangle=|\psi\rangle$ 

- It is enough to give the **generators** of *S*. For example:  $\{ZXXZI, IZXXZ, ZIZXX, XZIZX\}$
- Operators *O* that belong to the **normalizer** of *S*

 $O \in \mathbf{N}(\mathcal{S}) \qquad \iff \qquad O\mathcal{S} = \mathcal{S}O$ 

leave invariant the code space *C*. If they do not belong to the stabilizer, then they act non-trivially in the code subspace.

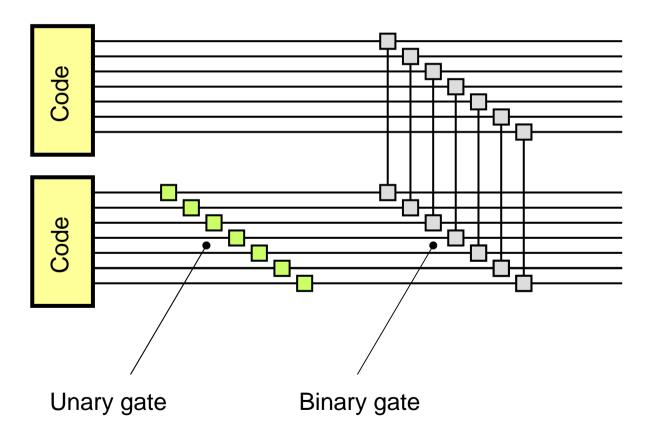
- A encoded state can be subject to **errors**.
- To correct them, we measure a set of generators of *S*. The results of the measurement compose the **syndrome** of the error. Errors can be corrected as long as the syndrome lets us distinguish among the possible errors.
- Since correctable errors always form a vector space, it is enough to consider Pauli operators, which form a basis.
- We say that a Pauli error *e* is **undetectable** if it belongs to N(*S*)-*S*. In such a case, the syndrome says nothing:

$$\forall s \in \mathcal{S} \qquad s \, e |\psi\rangle = e \, s' |\psi\rangle = e |\psi\rangle$$

• A set of Pauli errors *E* is correctable iff:

 $E^{\dagger}E \cap \mathbf{N}(\mathcal{S}) \subset \mathcal{S}$ 

Some stabilizer codes are specialy suitable for quantum computation. They allow to perform operations in a transversal and uniform way:



• Several codes allow the transversal implementation of

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix} \qquad K = \begin{pmatrix} 1 & 0\\ 0 & i \end{pmatrix} \qquad \Lambda = \begin{pmatrix} I_2 & 0\\ 0 & X \end{pmatrix}$$

which generate the **Clifford group**. This is useful for quantum information tasks such as teleportation or **entanglement distillation**.

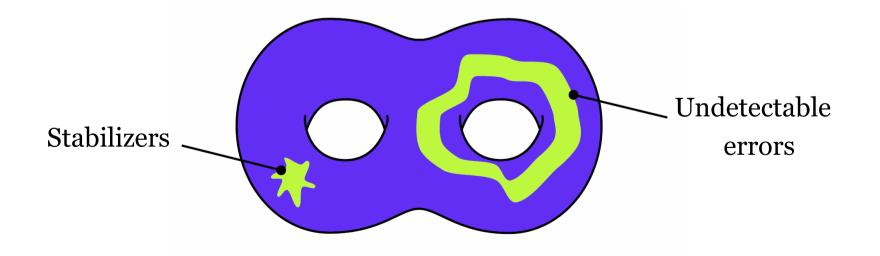
 Quantum Reed-Muller codes<sup>1</sup> are very special. They allow universal computation through transversal gates

$$K^{1/2} = \begin{pmatrix} 1 & 0 \\ 0 & i^{1/2} \end{pmatrix} \qquad \Lambda = \begin{pmatrix} I_2 & 0 \\ 0 & X \end{pmatrix}$$

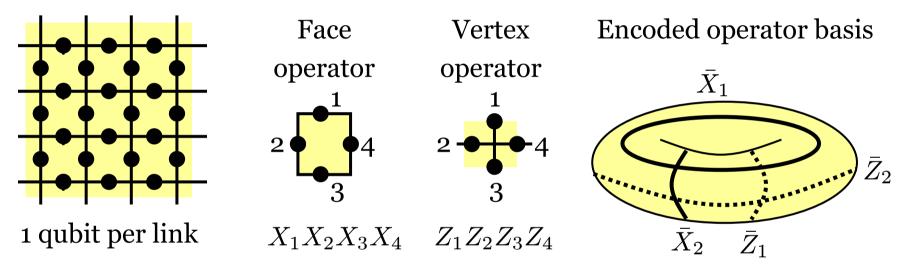
and transversal measurements of X and Z.

• We will see how both sets of operations can be transversally implemented in 2D and 3D topological color codes.

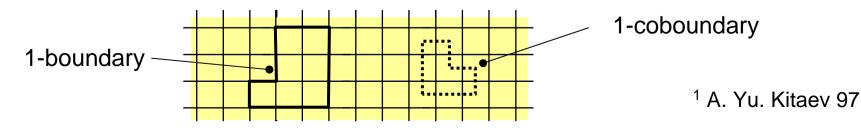
- In order to introduce the idea of a topological stabilizer code (TSC), we must consider a topological space in which our physical qubits are to be placed, for example a surface.
- A TSC is a stabilizer code in which the generators of the stabilizer are **local** and undetectable errors (or encoded operators) are **topologically nontrivial**.



• The first example of TSC were **surface codes**<sup>1</sup>, which are based on  $Z_2$  homology and cohomology.



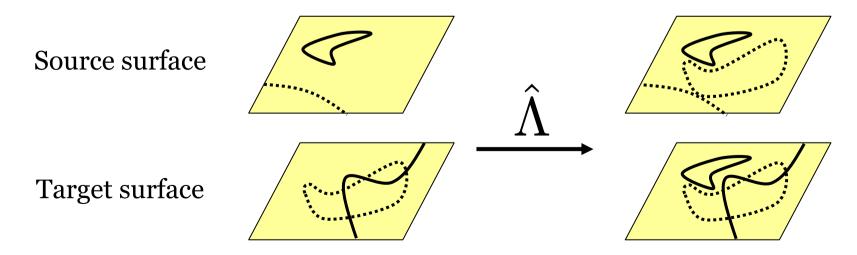
*S* gets identified with 1-boundaries and 1-coboundaries, and N(*S*) with 1-cycles and 1-cocycles.



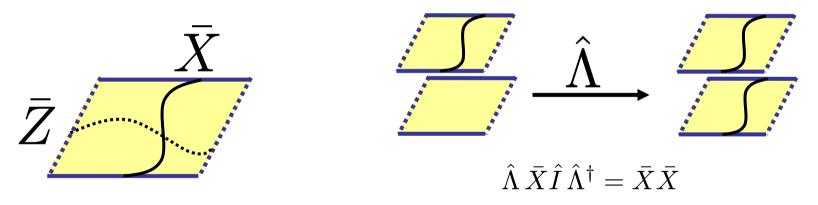
• The CNot gate can be implemented transversally on surface codes. First, its action under conjugation on operators is:

$$\Lambda: \begin{array}{ccc} {}_{IX} \longrightarrow {}_{IX} & {}_{IZ} \longrightarrow {}_{ZZ} \\ {}_{XI} \longrightarrow {}_{XX} & {}_{ZI} \longrightarrow {}_{ZI} \end{array}$$

• Thus the transversal action of the CNot on a surface code, at the level of operators, is simply to copy chains forward and cochains backwards.

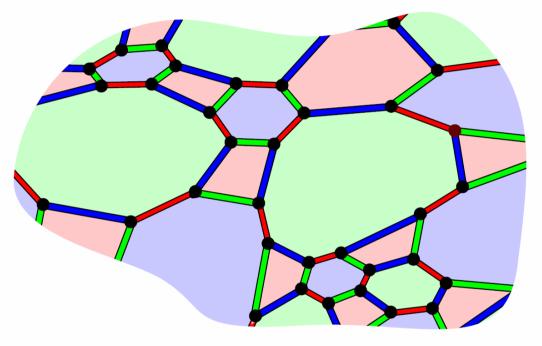


 Finally, to see the action of the tranversal CNOT on the code, we have to choose a Pauli basis for the encoded qubits. In the simplest example we have a single qubit in a square surface with suitable borders:



 Clearly the action of a transversal CNot is itself a CNot gate on the encoded qubits. However, this is the only gate we can get with surface codes. If we want to get further, we have to go beyond homology.

• A 2-colex is a **trivalent** 2-D lattice with **3-colored faces**.



- Edges can be 3-colored accordingly. Blue edges connect blue faces, and so on.
- The name 'colex' is for 'color complex'. *D*-colexes of arbitrary dimension can be defined. Their key feature is that the whole structure of the complex is contained in the 1-skeleton and the coloring of the edges.

• To construct a **color code** from a 2-colex, we place 1 qubit at each **vertex** of the lattice. The generators of *S* are **face operators**:

$$2 \sqrt[4]{5} B_{f}^{X} = X_{1}X_{2}X_{3}X_{4}X_{5}X_{6}$$

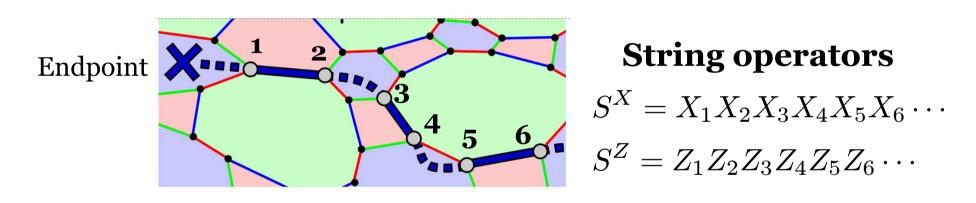
$$B_{f}^{Z} = Z_{1}Z_{2}Z_{3}Z_{4}Z_{5}Z_{6}$$

• Transversal Clifford gates should belong to **N**(*S*). We have:

$$\begin{split} \hat{H}B_f^X \hat{H}^{\dagger} &= B_f^Z & \hat{K}B_f^X \hat{K}^{\dagger} &= (-)^{\frac{v}{2}} B_f^X B_f^Z \\ \hat{H}B_f^Z \hat{H}^{\dagger} &= B_f^X & \hat{K}B_f^Z \hat{K}^{\dagger} &= B_f^Z \end{split}$$

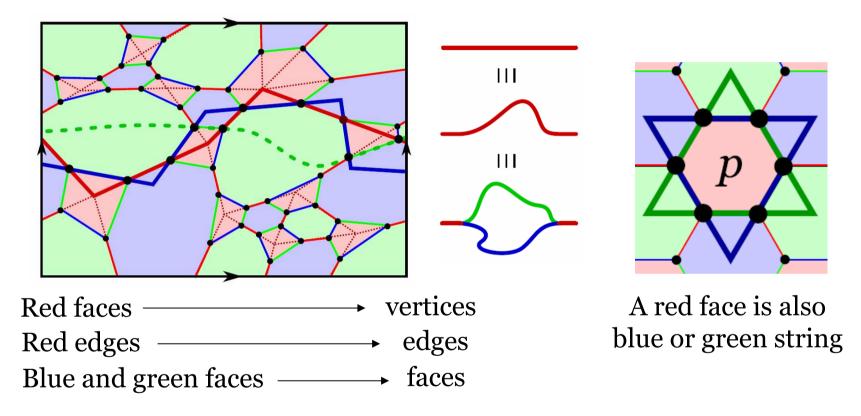
- Here v is the number of vertices in the face. If it is a multiple of 4 for every face, then K is in N(S). H always is.
- As for the CNot gate, it is clearly in N(S) (it is a CSS code).

- In order to understand 2-D color codes, we have to introduce string operators in the picture. As in surface codes, we play with  $Z_2$  homology. However, there is a new ingredient, color.
- A blue string is a collection of blue links:



Strings can have endpoints, located at faces of the same color. However, in that case the corresponding string and face operators will not commute. Therefore, a string operator belongs to N(S) iff the string has no endpoints.

• For each color we can form a **shrunk graph**. The red one is:



• Thus for each color homology works as in surface codes. The new feature is the possibility to **combine** homologous blue and red string operators of the same kind to get a green one.

- Since there are two independent colors, the number of encoded qubits should double that of a surface code. Lets check this for a surface **without boundary** using the Euler characteristic  $\chi = V + F E$  for any *shrunk* lattice.
- Face operators are subject to the **conditions**

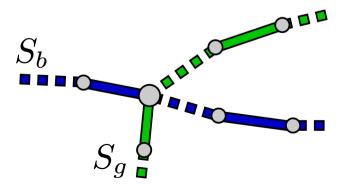
$$\prod_{f \in \bullet} B_f^{\sigma} = \prod_{f \in \bullet} B_f^{\sigma} = \prod_{f \in \bullet} B_f^{\sigma}$$

so that the total number of generators is g = 2(F + V - 2).

 The number of physical qubits is n = 2E. Therefore the number of encoded qubits q is twice the first Betti number of the manifold:

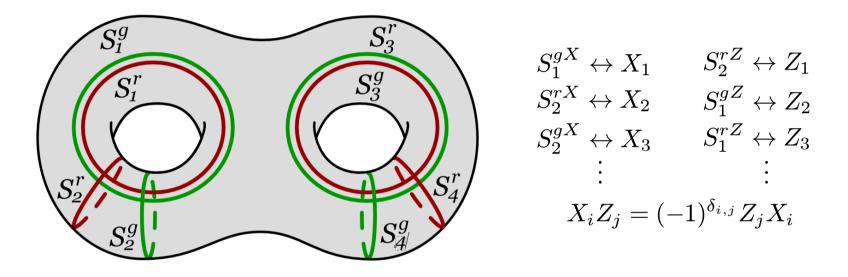
$$q = n - g = 4 - 2\chi = 2h_1$$

- In order to form a Pauli basis for the operators acting on encoded qubits, we can use as in surface codes those string operators (SO) that are not homologous to zero.
- To this end, we need the commutation rules for SO.
- Clearly SO of the same type (*X* or *Z*) always commute.
- A string is made up of edges with two vertices each. Therefore, two SO of the same color have an even number of qubits in common an they commute.
- SO of different colors can anticommute, but only if they cross an odd number of times:



$$\{S_b^X, S_g^Z\} = 0$$

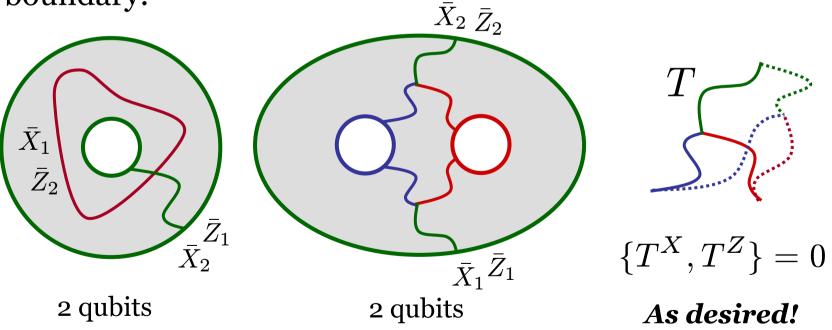
• Now we can construct the desired operator basis for the encoded qubits. In a 2-torus a possible choice is:



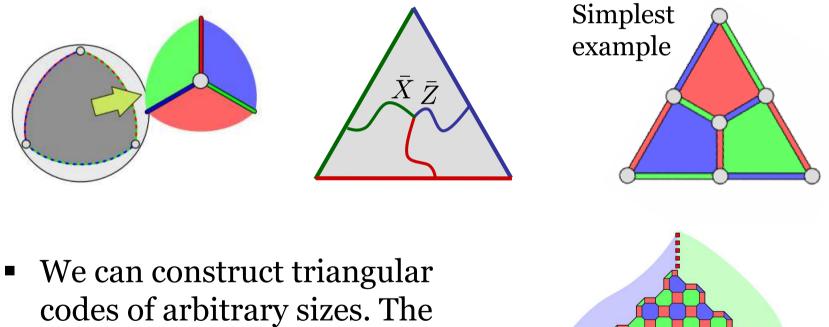
 However, if we apply the transversal *H* gate to such a code the resulting encoded gate is not *H*. The underlying reason is that for a string *S* we **never** have

$$\{S^X, S^Z\} = 0$$

- But we can consider surfaces with **boundary**. To this end, we take a sphere, which encodes no qubit, and **remove** faces.
- When a face is removed, the resulting boundary must have its color, and only strings of that color can end at the boundary.



• We can even encode a single qubit an remove the need for holes. If we remove a site and neighboring links and faces from a 2-colex in a sphere, we get a **triangular** code:



 We can construct triangular codes of arbitrary sizes. The vertices per face can be 4 and 8 so that *K* is in N(*S*).

• The transversal *H* clearly amounts to an encoded *H*:

$$H: \begin{array}{cc} X \longrightarrow Z \\ Z \longrightarrow X \end{array} \quad \hat{H}: \begin{array}{c} T^X \longrightarrow T^Z \\ T^Z \longrightarrow T^X \end{array} \qquad \begin{array}{c} T^X \end{array}$$

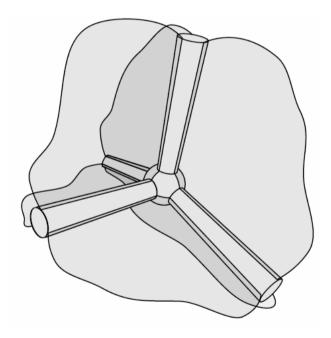
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• This is also true for *K*. The anticommutation properties of *T* imply that its support consists of an odd number of qubits:

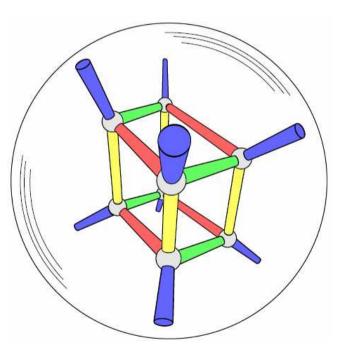
$$K: \begin{array}{cc} X \longrightarrow iXZ \\ Z \longrightarrow Z \end{array} \qquad \hat{K}: \begin{array}{c} T^X \longrightarrow \pm iT^X T^Z \\ T^Z \longrightarrow T^Z \end{array}$$

• Therefore, the **Clifford group** can be implemented transversally in triangular codes.

- **3-colexes** are tetravalent lattices with a particular local appearance such that their 3-cells can be 4-colored. They can be built in any compact 3-manifold without boundary.
- Edges can be colored accordingly, as in the 2-D case.

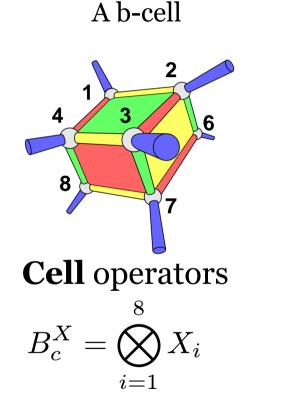


The neigborhood of a vertex.

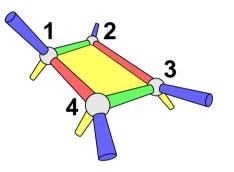


The simplest 3-colex in the projective space.

• This time the generators of *S* are face and (3-) cell operators.



A by-face separates band y-cells.



**Face** operators

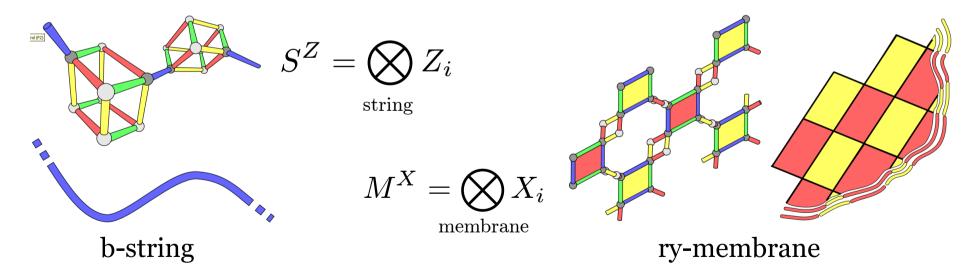
$$B_f^Z = \bigotimes_{i=1}^8 Z_i$$

 Therefore there are two different homology groups in the picture, those for 1-chains and for 2-chains. But in fact, due to Poincaré's duality they are the same.

- Strings are constructed as in 2-D, but now come in four colors. Branching is again possible.
- The new feature are membranes. They come in 6 color combinations and also have branching properties.

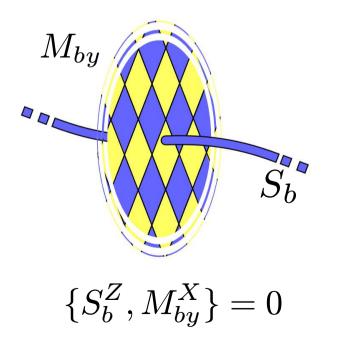
**String** operators

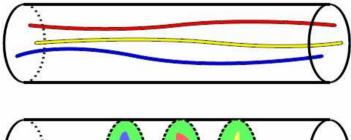
Membrane operators

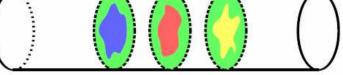


• There exist appropiate shunk complexes both for strings and for membranes.

- Now there are 3 indendent colors for strings (and similarly 3 color combinations for membranes). Therefore, we expect that the number of encoded qubits will be  $3h_1 = 3h_2$ .
- String and membrane operators always commute, unless they share a color and the string crosses an odd number of times the membrane.

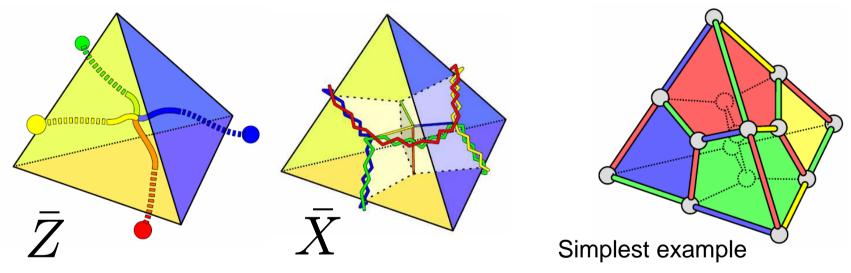






A pauli basis for the operators on the 3 qubits encoded in S<sup>2</sup>xS<sup>1</sup>.

- 3-Colexes cannot have a practical interest unless we allow boundaries. But this is just a matter of erasing cells. As in two dimensions, boundaries have the color of the erased cell.
- The analogue of triangular codes are tetrahedral codes, obtained by erasing a vertex from a 3-sphere.



• The desired transversal  $K^{1/2}$  gate can be implemented as long as faces have 4x vertices and cells 8x vertices.

# Conclusions

- *D*-colexes are *D*-valent complexes with certain coloring properties.
- Topological color codes are obtained from colexes. They have a richer structure than surface codes.
- 2-colexes allow the transversal implementation of Clifford operations.
- 3-colexes allow the transversal implementation of the same gates as Reed-Muller codes.