

## MP465 – Advanced Electromagnetism

### Problem Set 5

Due by 5pm on Friday, 8 May 2020

1. A centre-fed linear antenna consists of two conducting wires, each of length  $L$ , arranged horizontally such that an alternating current is fed into the (insulated) gap between them. If we take the antenna to be along the  $y$ -axis, then the current density is given by

$$\vec{J}(t, x, y, z) = \begin{cases} I_0 \left(1 - \frac{|y|}{L}\right) \delta(x)\delta(z)\hat{e}_y \sin \omega t & \text{for } -L \leq y \leq L, \\ \vec{0} & \text{otherwise,} \end{cases}$$

where  $I_0$  is a real constant.

- (a) Find the complex electric dipole amplitude associated with this current, namely, the vector  $\vec{p}_0$  such that  $\vec{p}(t) = \text{Re}[\vec{p}_0 e^{-i\omega t}]$ .
  - (b) Find the time-averaged power distribution  $d\bar{P}/d\Omega$  as a function of the spherical angles  $\theta$  and  $\phi$ , and use it to compute the total time-averaged radiated power  $\bar{P}$ .
2. If we do *not* assume that the time-dependence of our sources is periodic, then the far-zone approximation for the magnetic field is

$$\vec{B}(t, \vec{r}) \approx \frac{\mu_0}{4\pi c} \frac{\ddot{\vec{p}}(t - r/c) \times \hat{e}_r}{r},$$

where

$$\vec{p}(t) = \int \rho(t, \vec{r}) \vec{r} d^3\vec{r}.$$

is the time-dependent electric dipole moment of the sources. (The electric field is still given by  $\vec{E} \approx c\vec{B} \times \hat{e}_r$ .)

- (a) Without assuming anything about the way the charge density depends on time, show from the continuity equation that the derivative of  $\vec{p}(t)$  is

$$\dot{\vec{p}}(t) = \int \vec{J}(t, \vec{r}) d^3\vec{r}.$$

- (b) A point charge  $q$  undergoes a constant acceleration  $\vec{a} = a\hat{e}_z$ . Find the total power radiated by this charge.
3. We know from lecture that when we have an inertial frame  $\mathcal{S}'$  moving at a constant velocity  $\vec{v}$  relative to an inertial frame  $\mathcal{S}$ , then the electric and magnetic fields measured in  $\mathcal{S}'$  are related to the ones in  $\mathcal{S}$  by

$$\begin{aligned} E'_{\parallel} &= E_{\parallel}, & \vec{E}'_{\perp} &= \gamma(v) \left( \vec{E}_{\perp} + \vec{v} \times \vec{B} \right), \\ B'_{\parallel} &= B_{\parallel}, & \vec{B}'_{\perp} &= \gamma(v) \left( \vec{B}_{\perp} - \frac{\vec{v}}{c^2} \times \vec{E} \right), \end{aligned}$$

where  $\parallel$  and  $\perp$  denote, respectively, the components parallel and perpendicular to  $\vec{v}$ .

In a particular frame  $\mathcal{S}$ , we have an electromagnetic field given by

$$\vec{E} = E_0 \hat{e}_y, \quad \vec{B} = \frac{\sqrt{2}E_0}{c} (\hat{e}_x - \hat{e}_y),$$

where  $E_0$  is a constant. Consider the frame  $\mathcal{S}'$  in which the electric and magnetic fields are antiparallel to each other; if  $v_x = 0$ , find  $v_y$  and  $v_z$ . (Hint: if  $\vec{E}'$  and  $\vec{B}'$  are antiparallel, this means that there is some positive constant  $\lambda$  such that  $\vec{B}' = -\lambda \vec{E}'/c$ .)

4. The field strength and dual field strength of an electromagnetic field described by the 4-potential  $A^\mu = (\Phi/c, \vec{A})$  are, respectively,

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu, \quad \star F^{\mu\nu} = \frac{1}{2} \epsilon^{\mu\nu\lambda\rho} F_{\lambda\rho}.$$

Show that

$$\star F^{\mu\nu} F_{\mu\nu} = -\frac{4}{c} \vec{E} \cdot \vec{B},$$

and explain why this implies that  $\vec{E} \cdot \vec{B}$  is Lorentz-invariant.

## VECTOR CALCULUS FORMULAE

1. Cartesian coordinates  $(x, y, z)$  with constant unit direction vectors  $\hat{e}_x$ ,  $\hat{e}_y$ ,  $\hat{e}_z$

- position vector:  $\vec{r} = x\hat{e}_x + y\hat{e}_y + z\hat{e}_z$
- line element:  $d\vec{r} = dx\hat{e}_x + dy\hat{e}_y + dz\hat{e}_z$   
 surface element:  $d\vec{\sigma} = dy\,dz\hat{e}_x + dx\,dz\hat{e}_y + dx\,dy\hat{e}_z$   
 volume element:  $d^3\vec{r} = dx\,dy\,dz$
- gradient of a function  $f(x, y, z)$ :

$$\vec{\nabla}f = \frac{\partial f}{\partial x}\hat{e}_x + \frac{\partial f}{\partial y}\hat{e}_y + \frac{\partial f}{\partial z}\hat{e}_z$$

- divergence of a vector  $\vec{A}(x, y, z) = A_x(x, y, z)\hat{e}_x + A_y(x, y, z)\hat{e}_y + A_z(x, y, z)\hat{e}_z$ :

$$\vec{\nabla} \cdot \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$

- curl of a vector  $\vec{A}(x, y, z) = A_x(x, y, z)\hat{e}_x + A_y(x, y, z)\hat{e}_y + A_z(x, y, z)\hat{e}_z$ :

$$\vec{\nabla} \times \vec{A} = \left( \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{e}_x + \left( \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{e}_y + \left( \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{e}_z$$

- Laplacian of a function  $f(x, y, z)$ :

$$\nabla^2 f = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$$

2. Cylindrical coordinates  $(r, \phi, z)$  with unit direction vectors  $\hat{e}_r$ ,  $\hat{e}_\phi$ ,  $\hat{e}_z$

- relation to Cartesian coordinates:  $x = r \cos \phi$ ,  $y = r \sin \phi$ ,  $z$  unchanged
- relation to Cartesian unit vectors:

$$\left. \begin{array}{l} \hat{e}_r = \cos \phi \hat{e}_x + \sin \phi \hat{e}_y \\ \hat{e}_\phi = -\sin \phi \hat{e}_x + \cos \phi \hat{e}_y \end{array} \right\} \leftrightarrow \left\{ \begin{array}{l} \hat{e}_x = \cos \phi \hat{e}_r - \sin \phi \hat{e}_\phi \\ \hat{e}_y = \sin \phi \hat{e}_r + \cos \phi \hat{e}_\phi \end{array} \right.$$

with  $\hat{e}_z$  the same for both systems.

- position vector:  $\vec{r} = r\hat{e}_r + z\hat{e}_z$
- line element:  $d\vec{r} = dr\hat{e}_r + r d\phi\hat{e}_\phi + dz\hat{e}_z$   
 surface element:  $d\vec{\sigma} = r d\phi dz\hat{e}_r + dr dz\hat{e}_\phi + r dr d\phi\hat{e}_z$   
 volume element:  $d^3\vec{r} = r dr d\phi dz$
- gradient of a function  $f(r, \phi, z)$ :

$$\vec{\nabla} f = \frac{\partial f}{\partial r}\hat{e}_r + \frac{1}{r}\frac{\partial f}{\partial \phi}\hat{e}_\phi + \frac{\partial f}{\partial z}\hat{e}_z$$

- divergence of a vector  $\vec{A}(r, \phi, z) = A_r(r, \phi, z)\hat{e}_r + A_\phi(r, \phi, z)\hat{e}_\phi + A_z(r, \phi, z)\hat{e}_z$ :

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{r}\frac{\partial}{\partial r}(rA_r) + \frac{1}{r}\frac{\partial A_\phi}{\partial \phi} + \frac{\partial A_z}{\partial z}$$

- curl of a vector  $\vec{A}(r, \phi, z) = A_r(r, \phi, z)\hat{e}_r + A_\phi(r, \phi, z)\hat{e}_\phi + A_z(r, \phi, z)\hat{e}_z$ :

$$\vec{\nabla} \times \vec{A} = \left(\frac{1}{r}\frac{\partial A_z}{\partial \phi} - \frac{\partial A_\phi}{\partial z}\right)\hat{e}_r + \left(\frac{\partial A_r}{\partial z} - \frac{\partial A_z}{\partial r}\right)\hat{e}_\phi + \frac{1}{r}\left(\frac{\partial}{\partial r}(rA_\phi) - \frac{\partial A_r}{\partial \phi}\right)\hat{e}_z$$

- Laplacian of a function  $f(r, \phi, z)$ :

$$\nabla^2 f = \frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial f}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 f}{\partial \phi^2} + \frac{\partial^2 f}{\partial z^2}$$

3. Spherical coordinates  $(r, \theta, \phi)$  with unit direction vectors  $\hat{e}_r, \hat{e}_\theta, \hat{e}_\phi$

- relation to Cartesian coordinates:  $x = r \sin \theta \cos \phi, y = r \sin \theta \sin \phi, z = r \cos \theta$
- relation to Cartesian unit vectors:

$$\begin{aligned} & \left. \begin{aligned} \hat{e}_r &= \sin \theta \cos \phi \hat{e}_x + \sin \theta \sin \phi \hat{e}_y + \cos \theta \hat{e}_z \\ \hat{e}_\theta &= \cos \theta \cos \phi \hat{e}_x + \cos \theta \sin \phi \hat{e}_y - \sin \theta \hat{e}_z \\ \hat{e}_\phi &= -\sin \phi \hat{e}_x + \cos \phi \hat{e}_y \end{aligned} \right\} \\ \Leftrightarrow & \left\{ \begin{aligned} \hat{e}_x &= \sin \theta \cos \phi \hat{e}_r + \cos \theta \cos \phi \hat{e}_\theta - \sin \phi \hat{e}_\phi \\ \hat{e}_y &= \sin \theta \sin \phi \hat{e}_r + \cos \theta \sin \phi \hat{e}_\theta + \cos \phi \hat{e}_\phi \\ \hat{e}_z &= \cos \theta \hat{e}_r - \sin \theta \hat{e}_\theta \end{aligned} \right. \end{aligned}$$

- position vector:  $\vec{r} = r \hat{e}_r$
- line element:  $d\vec{r} = dr \hat{e}_r + r d\theta \hat{e}_\theta + r \sin \theta d\phi \hat{e}_\phi$   
surface element:  $d\vec{\sigma} = r^2 \sin \theta d\theta d\phi \hat{e}_r + r \sin \theta dr d\phi \hat{e}_\theta + r dr d\theta \hat{e}_\phi$   
volume element:  $d^3\vec{r} = r^2 \sin \theta dr d\theta d\phi$
- gradient of a function  $f(r, \theta, \phi)$ :

$$\vec{\nabla} f = \frac{\partial f}{\partial r} \hat{e}_r + \frac{1}{r} \frac{\partial f}{\partial \theta} \hat{e}_\theta + \frac{1}{r \sin \theta} \frac{\partial f}{\partial \phi} \hat{e}_\phi$$

- divergence of a vector  $\vec{A}(r, \theta, \phi) = A_r(r, \theta, \phi) \hat{e}_r + A_\theta(r, \theta, \phi) \hat{e}_\theta + A_\phi(r, \theta, \phi) \hat{e}_\phi$ :

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta A_\theta) + \frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi}$$

- curl of a vector  $\vec{A}(r, \theta, \phi) = A_r(r, \theta, \phi) \hat{e}_r + A_\theta(r, \theta, \phi) \hat{e}_\theta + A_\phi(r, \theta, \phi) \hat{e}_\phi$ :

$$\begin{aligned} \vec{\nabla} \times \vec{A} &= \frac{1}{r \sin \theta} \left( \frac{\partial}{\partial \theta} (\sin \theta A_\phi) - \frac{\partial A_\theta}{\partial \phi} \right) \hat{e}_r + \left( \frac{1}{r \sin \theta} \frac{\partial A_r}{\partial \phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_\phi) \right) \hat{e}_\theta \\ &+ \frac{1}{r} \left( \frac{\partial}{\partial r} (r A_\theta) - \frac{\partial A_r}{\partial \theta} \right) \hat{e}_\phi \end{aligned}$$

- Laplacian of a function  $f(r, \theta, \phi)$ :

$$\nabla^2 f = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial f}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial f}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 f}{\partial \phi^2}$$