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THE NATIONAL UNIVERSITY OF IRELAND MAYNOOTH

**MATHEMATICAL PHYSICS**

**REPEAT EXAMINATION**

**2018–2019**

**Computational Physics 2**

**MP468C**

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**Time allowed: 2 hours**

**Answer ALL questions**

1. (a) Use the transformation method to construct a recipe for obtaining pseudo-random numbers with probability distribution

$$p(y) = \frac{3y^2}{y^3 + 1}, \quad y \in [0, \sqrt[3]{e-1}]. \quad (1.1)$$

You may assume a generator of uniform pseudo-random numbers between 0 and 1 is given

**[11 marks]**

- (b) List the main steps of the rejection method for generating a pseudo-random number distributed according to  $f(x)$ , given a constant  $M \in \mathbb{R}$  and generators for uniform pseudo-random numbers between 0 and 1 and pseudo-random numbers distributed according to  $g(x)$  such that  $f(x) < Mg(x)$ , for all  $x \in \mathbb{R}$ . Prove that the rejection method produces a variable  $Y$  distributed according to  $f(x)$ .

**[15 marks]**

- (c) If  $x_i$  are  $N$  independent uniformly distributed random points within a  $d$ -dimensional volume  $V$ , and

$$I_{MC} = \frac{V}{N} \sum_{i=1}^N f(x_i), \quad (1.2)$$

show that the expectation value of  $I_{MC}$  is equal to the integral of the function  $f$  over the volume  $V$ ,

$$\langle I_{MC} \rangle = \int_V f(x) dx. \quad (1.3)$$

Explain how this relation can be used to compute the integral  $I$  using Monte Carlo integration.

**[12 marks]**

- (d) Assuming you have a random number generator to generate pseudo-random numbers  $x$  with the following distribution:

$$p(x) = e^{-x}, \quad x \in [0, \infty). \quad (1.4)$$

explain how you would use Monte Carlo integration with importance sampling to compute the following integral:

$$I = \int_0^{\infty} 4xe^{(4-x)} dx \quad (1.5)$$

**[12 marks]**

2. (a) Using symmetric first and second derivatives, write down the discretised version of the equation

$$A \frac{\partial^2 \phi}{\partial x^2} + B \frac{\partial^2 \phi}{\partial y^2} + C \frac{\partial \phi}{\partial x} + D \frac{\partial \phi}{\partial y} = \rho(x, y), \quad 0 \leq x, y \leq L, \quad (2.1)$$

where  $A, B, C$  and  $D$  are known constants and  $\rho(x, y)$  is a known function of  $x$  and  $y$ , on a square symmetric grid of  $N \times N$  points with zero Dirichlet boundary conditions.

[13 marks]

- (b) Explain how the resulting equation can be written as a matrix equation,  $M\Phi = B$ , where  $M$  is a sparse  $N^2 \times N^2$  matrix and  $B$  is a known vector of length  $N^2$ . Write down expressions for  $M$  and  $B$ , taking the boundary conditions into account.

[10 marks]

- (c) Write down the Forward Time Centred Space discretisation scheme for the equation below, assuming equal lattice spacings  $\delta x = \delta y = a$  in both space directions, and a spacing  $\Delta t$  in the time direction.

$$\frac{\partial \phi}{\partial t} = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} - \rho(x, y), \quad (2.2)$$

[12 marks]

- (d) Using the von Neumann stability criterion for this scheme,  $\Delta t \leq a^2/4$ , derive the *Jacobi* method for solving the Poisson equation, and explain how it may be modified to obtain the *Gauss-Seidel* method.

[10 marks]

- (e) Assuming that each iteration reduces the difference between your estimate and the true solution by a factor  $\rho_s$  (called the spectral radius), find how many iterations are required to reduce this difference by a factor  $10^{-p}$ .

For the Poisson equation on a square  $N \times N$  grid with homogeneous Dirichlet boundary conditions, the spectral radii for the Jacobi and Gauss-Seidel methods are given by

$$\text{Jacobi: } \rho_J = \cos\left(\frac{\pi}{N}\right), \quad \text{Gauss-Seidel: } \rho_{GS} = \cos^2\left(\frac{\pi}{N}\right). \quad (2.3)$$

Use this to show that the Gauss-Seidel method converges twice as fast as the Jacobi method, and that the number of iterations required for both to converge increases as  $N^2$  in the limit of large  $N$ .

[5 marks]

## Solutions: Question 1

- (a) We want to find a function  $f$  such that, given  $X$  is uniformly distributed between 0 and 1,  $Y = f(X)$  is distributed according to

$$p(y) = \frac{3y^2}{y^3 + 1}, \quad (0.4)$$

We know that

$$\text{Probability } X \in [a, b] = \text{Probability } Y \in [f(a), f(b)], \quad (0.5)$$

$$\implies \int_a^b P_X(x) dx = \int_{f(a)}^{f(b)} P_Y(y) dy, \quad (0.6)$$

$$\implies \int_0^x dx' = \int_0^y \frac{3y'^2}{y'^3 + 1} dy', \quad (0.7)$$

$$\implies x = \int_1^{y^3+1} \frac{d\alpha}{\alpha} = \ln(y^3 + 1) \quad (0.8)$$

Inverting this gives

$$y = \sqrt[3]{e^x - 1}. \quad (0.9)$$

Hence, generating a uniformly distributed number  $X$  between 0 and 1 and applying the function  $f(x) = \sqrt[3]{e^x - 1}$  produces a number  $Y$  which is distributed according to (0.4).

[11 marks]

- (b) Given a constant  $M \in \mathbb{R}$ , the following three steps will produce a random number  $Y$  distributed according to  $f(x)$  given generators for producing random numbers  $u$  distributed uniformly between 0 and 1 and  $X$  distributed according to  $g(x)$ .

- Generate a random number  $X$  according to  $f(x)$ .
- Generate a uniformly distributed random number  $u$  between 0 and 1.
- If  $u < f(X)/Mg(X)$ , accept  $Y = X$ . Other wise reject  $X$  and execute these three steps again.

To prove  $Y$  is distributed according to  $f(x)$ , we first show that the probability of  $Y$  being less than  $x$  is given by

$$P(Y < x) = \int_{-\infty}^x f(\tilde{x}) d\tilde{x}. \quad (0.10)$$

We note, for  $Y$  to be less than  $x$ , two things must be true. Firstly, the random number  $u$  must be less than  $f(x)/Mg(x)$ . Then, provided that's true,  $X$  must be less than  $x$ . Hence, we have

$$P(Y < x) = P(X < x | u < f(x)/Mg(x)), \quad (0.11)$$

$$= \frac{P(X < x, u < f(x)/Mg(x))}{u < f(x)/Mg(x)}. \quad (0.12)$$

We now note, that since  $X$  and  $u$  are independent random variables, the tuple  $(X, u)$  is distributed in the plane according to the product of distributions for  $X$  and  $u$ .

$$(X, u) \sim P(x, y) = g(x)P_{\text{uni}}^{[0,1]}(y). \quad (0.13)$$

Rewriting the probabilities appearing in (0.12) as integrals of the above distribution yields

$$P(Y < x) = \frac{\int_{-\infty}^x \left( \int_0^{f(\tilde{x})/Mg(\tilde{x})} g(\tilde{x}) du \right) d\tilde{x}}{\int_{-\infty}^{+\infty} \left( \int_0^{f(\tilde{x})/Mg(\tilde{x})} g(\tilde{x}) du \right) d\tilde{x}} \quad (0.14)$$

$$= \frac{\int_{-\infty}^x [f(\tilde{x})/Mg(\tilde{x})] g(\tilde{x}) d\tilde{x}}{\int_{-\infty}^{+\infty} [f(\tilde{x})/Mg(\tilde{x})] g(\tilde{x}) d\tilde{x}} \quad (0.15)$$

$$= \frac{\int_{-\infty}^x f(\tilde{x}) d\tilde{x}}{\int_{-\infty}^{+\infty} f(\tilde{x}) d\tilde{x}} \quad (0.16)$$

$$= \int_{-\infty}^x f(\tilde{x}) d\tilde{x} \quad (0.17)$$

The probability density of  $Y$  is given by the derivative of its cumulative distribution. So the distribution of  $Y$  must be

$$P_Y(y) = \frac{d}{dx} \left( \int_{-\infty}^x f(\tilde{x}) d\tilde{x} \right) \Big|_{x=y} = f(y), \quad (0.18)$$

proving that  $Y$  is distributed according to  $f(x)$ .

**[15 marks]**

- (c) If  $X$  is uniformly distributed inside the  $d$ -dimensional volume  $V$ , then the expectation value of the quantity  $f(X)$  is

$$\langle f \rangle = \int_V f(x) P_{\text{uni}}^V(x) dx = \int_V \frac{f(x)}{V} dx. \quad (0.19)$$

Hence the expectation value of  $I_{\text{MC}}$  is

$$\langle I_{\text{MC}} \rangle = \frac{V}{N} \sum_{i=1}^N \langle f \rangle = \frac{V}{N} (N \langle f \rangle), \quad (0.20)$$

$$= V \int_V \frac{f(x)}{V} dx = \int_V f(x) dx. \quad (0.21)$$

For the first equality we used the fact that the  $x_i$ s are independent. Hence the expectation value of  $I_{\text{MC}}$  is the integral of  $f$  over the region  $V$ .

[12 marks]

(d) Rewriting the integrand yields

$$I = \int_0^\infty 4xe^{(4-x)} dx \quad (0.22)$$

$$= 4e^4 \int_0^\infty xe^{-x} dx \quad (0.23)$$

Hence, to compute  $I$  using Monte Carlo integration with importance sampling one can generate  $N$  pseudo-random numbers  $x_i$  distributed under  $p(x)$  and compute

$$I = \frac{4e^4}{N} \sum_{i=1}^N x \quad (0.24)$$

[12 marks]

## Solutions: Question 2

(a) The symmetric finite difference equation for the first derivative of a function  $f$  is

$$f'(x) = \frac{f(x+a) - f(x-a)}{2a}. \quad (0.25)$$

The symmetric finite difference equation for the second derivative of a function  $f$  is

$$f''(x) = \frac{f(x+a) - 2f(x) + f(x-a)}{a^2}. \quad (0.26)$$

We discretise the defined square region of the  $x, y$ -plane into a symmetric  $(N + 2) \times (N + 2)$  grid, with lattice spacing  $a = \frac{L}{N+1}$ . For a function  $\phi(x, y)$  on the interior points of the lattice we write

$$\phi(x, y) = \phi(ia, ja) = \phi_{i,j}, \quad (0.27)$$

where  $i, j = 1 \dots N$ . The first and second derivatives, at interior points of the lattice, are then given by

$$\frac{\partial \phi}{\partial x}(x, y) = \frac{\phi_{i+1,j} - \phi_{i-1,j}}{2a}, \quad (0.28)$$

$$\frac{\partial \phi}{\partial y}(x, y) = \frac{\phi_{i,j+1} - \phi_{i,j-1}}{2a}, \quad (0.29)$$

$$\frac{\partial^2 \phi}{\partial x^2}(x, y) = \frac{\phi_{i+1,j} - 2\phi_{i,j} + \phi_{i-1,j}}{a^2}, \quad (0.30)$$

$$\frac{\partial^2 \phi}{\partial y^2}(x, y) = \frac{\phi_{i,j+1} - 2\phi_{i,j} + \phi_{i,j-1}}{a^2}. \quad (0.31)$$

Substituting these into the differential equation yields

$$\rho(x, y) = A \frac{\partial^2 \phi}{\partial x^2} + B \frac{\partial^2 \phi}{\partial y^2} + C \frac{\partial \phi}{\partial x} + D \frac{\partial \phi}{\partial y}, \quad (0.32)$$

$$= A \frac{\phi_{i+1,j} - 2\phi_{i,j} + \phi_{i-1,j}}{a^2} + B \frac{\phi_{i,j+1} - 2\phi_{i,j} + \phi_{i,j-1}}{a^2} \quad (0.33)$$

$$+ C \frac{\phi_{i+1,j} - \phi_{i-1,j}}{2a} + D \frac{\phi_{i,j+1} - \phi_{i,j-1}}{2a}, \quad (0.34)$$

$$= \frac{1}{a^2} \left[ \left( A + \frac{Ca}{2} \right) \phi_{i+1,j} + \left( A - \frac{Ca}{2} \right) \phi_{i-1,j} \right. \quad (0.35)$$

$$\left. + \left( B + \frac{Da}{2} \right) \phi_{i,j+1} + \left( B - \frac{Da}{2} \right) \phi_{i,j-1} - 4\phi_{i,j} \right] \quad (0.36)$$

[13 marks]

(b) We first define  $\tilde{\rho}_{i,j} = a^2 \rho(x, y)$  so we can write

$$\tilde{\rho}_{i,j} = \left( A + \frac{Ca}{2} \right) \phi_{i+1,j} + \left( A - \frac{Ca}{2} \right) \phi_{i-1,j} \quad (0.37)$$

$$+ \left( B + \frac{Da}{2} \right) \phi_{i,j+1} + \left( B - \frac{Da}{2} \right) \phi_{i,j-1} - 4\phi_{i,j}. \quad (0.38)$$

To write this as a matrix equation we number the sites of the  $N \times N$  lattice 1 to  $N^2$ . The number we assign to the site  $(i, j)$  is  $n = i + Nj$ . Then we can list the values of  $\phi_{i,j}$  and  $\tilde{\rho}_{i,j}$ , from 1 to  $N^2$ , in column vectors  $\Phi$  and  $B$  respectively. The above equation then turns into the matrix equation  $M\Phi = B$ , where  $M$  is a  $N^2 \times N^2$  matrix whose components are given by

$$M_{m,n} = \left( A + \frac{Ca}{2} \right) \delta_{m,n+1} + \left( A - \frac{Ca}{2} \right) \delta_{m,n-1} \quad (0.39)$$

$$+ \left( B + \frac{Da}{2} \right) \delta_{m,n+N} + \left( B - \frac{Da}{2} \right) \delta_{m,n-N} - 4\delta_{m,n} \quad (0.40)$$

and the vector  $B$  is given by

$$B_n = \tilde{\rho}_n = \tilde{\rho}_{i,j}, \quad (0.41)$$

where  $n = i + Nj$ .

Equations involving boundary terms, in the set of linear equations  $M\Phi = B$ , are treated differently since  $\phi_{0,j} = \phi_{N+1,j} = \phi_{i,0} = \phi_{i,N+1} = 0$  (zero Dirichlet boundary conditions). This amounts to terms appearing in 0.40 being set to zero for certain values of  $m$  (certain row equations). Namely, the following terms in 0.40 are set to zero, (here  $i(m)$  and  $j(m)$  are the original indices.)

$$\begin{aligned} \text{When } i(m) = 1, \quad & \delta_{m,n-1} = 0, \\ \text{When } i(m) = N, \quad & \delta_{m,n+1} = 0, \\ \text{When } j(m) = 1, \quad & \delta_{m,n-N} = 0, \\ \text{When } j(m) = N, \quad & \delta_{m,n+N} = 0. \end{aligned}$$

In general, to implement Dirichlet boundary conditions, one must subtract the boundary values from  $\tilde{\rho}_n$  appropriately to form the vector  $B$ . However, since we want to implement zero Dirichlet boundary conditions, this amounts to subtracting zero, leaving the equation 0.41 unaltered.

**[10 marks]**

- (c) If we use a lattice with equal lattice spacings  $\delta x = \delta y = a$  in both space directions, and a spacing  $\Delta t$  in the time direction, the Forward Time Centred Space discretisation scheme for this equation is the following finite difference equation.

$$\frac{\partial \phi}{\partial t} = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} - \rho(x, y), \quad (0.42)$$

$$\frac{\phi_{i,j}^{n+1} - \phi_{i,j}^n}{\Delta t} = \frac{\phi_{i+1,j}^n - 2\phi_{i,j}^n + \phi_{i-1,j}^n}{a^2} + \frac{\phi_{i,j+1}^n - 2\phi_{i,j}^n + \phi_{i,j-1}^n}{a^2} - \rho_{i,j} \quad (0.43)$$

$$\phi_{i,j}^{n+1} = \left(1 - \frac{4\Delta t}{a^2}\right) \phi_{i,j}^n + \frac{\Delta t}{a^2} (\phi_{i+1,j}^n + \phi_{i-1,j}^n + \phi_{i,j+1}^n + \phi_{i,j-1}^n) - \Delta t \rho_{i,j} \quad (0.44)$$

[12 marks]

- (d) The von Neumann stability criterion for this scheme is  $\Delta t \leq a^2/4$ . The Jacobi method for solving the Poisson equation is to use the largest possible time step size  $\Delta t = a^2/4$ . This amounts to iterating the following equation to evolve an arbitrary initial state  $\phi$  until it converges to a stationary state.

$$\phi_{i,j}^{n+1} = \frac{1}{4} (\phi_{i+1,j}^n + \phi_{i-1,j}^n + \phi_{i,j+1}^n + \phi_{i,j-1}^n) - \frac{a^2}{4} \rho_{i,j}. \quad (0.45)$$

The above procedure can be modified to obtain the GaussSeidel method by using values of  $\phi(n+1)$  that have already been computed to calculate each  $\phi_{i,j}^{n+1}$ . This amounts to using the following equation

$$\phi_{i,j}^{n+1} = \frac{1}{4} (\phi_{i+1,j}^n + \phi_{i-1,j}^{n+1} + \phi_{i,j+1}^n + \phi_{i,j-1}^{n+1}) - \frac{a^2}{4} \rho_{i,j}. \quad (0.46)$$

[10 marks]

- (e) Assuming that each iteration reduces the difference between the estimate and the true solution by a factor  $\rho_s$ , the number of iterations  $n$  required to reduce this difference by a factor  $10^{-p}$  is given by

$$\rho_s^n = 10^{-p}, \quad (0.47)$$

$$\implies n \ln(\rho_s) = -p \ln(10), \quad (0.48)$$

$$\implies n = \frac{-p \ln(10)}{\ln(\rho_s)}. \quad (0.49)$$

For the Jacobi method we have  $\rho_s = \rho_J = \cos(\frac{\pi}{N})$ . So the number of iterations of the Jacobi method needed to reduce the difference by a factor of  $10^{-p}$  is

$$n_J = \frac{-p \ln(10)}{\ln(\rho_J)} = \frac{-p \ln(10)}{\ln(\cos(\frac{\pi}{N}))}. \quad (0.50)$$

For the GaussSeidel method we have  $\rho_s = \rho_{GS} = \cos^2(\frac{\pi}{N})$ . So the number of iterations of the GaussSeidel method needed to reduce the difference by a factor of  $10^{-p}$  is

$$n_{GS} = \frac{-p \ln(10)}{\ln(\rho_{GS})} = \frac{-p \ln(10)}{\ln(\cos^2(\frac{\pi}{N}))} = \frac{-p \ln(10)}{2 \ln(\cos(\frac{\pi}{N}))} = \frac{n_J}{2} \quad (0.51)$$

Hence the GaussSeidel converges twice as fast as the Jacobi method.

**[5 marks]**